

# Heat Transfer

## Computational Laboratories

### One-Dimensional, Transient Conduction (Laboratory II)

Space- and time-dependent conduction heat transfer in large plane walls, long rods, and spheres initiated by convection heat transfer across its boundaries

# Transient Conduction - Introduction

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- A transient (unsteady or time-dependent) heat conduction process is initiated whenever a thermal equilibrium state of a system is perturbed.
- A perturbation on a system thermal equilibrium state can be induced by a change in:
  - surface convection conditions ( $T_\infty$  or  $h$ );
  - surface radiation conditions ( $T_{\text{SUR}}$  or  $h_r$ );
  - surface heat flux ( $q_s''$ ) or surface temperature ( $T_s$ ); and
  - internal energy generation ( $\dot{q}$ ).
- Transient heat conduction processes can be modelled through analytical or numerical means:
  - Lumped system analysis (overall energy balance);
  - Exact solutions to the heat diffusion equation; and
  - Finite difference, finite element or finite volume methods.

# Transient Conduction - Temperature Gradients

## Importance of Solid Temperature Spatial Resolution

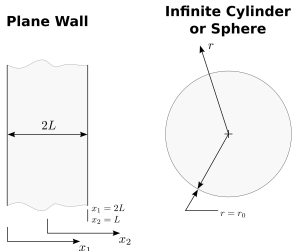
- For a transient conduction process in a solid driven by convection heat transfer across its boundaries, Biot number ( $Bi$ ) determines if the spatial thermal gradients within the solid are negligible or not.

$$Bi = \frac{hL_c}{k} = \frac{\text{Conduction resistance within the solid}}{\text{Convection resistance between the solid and the fluid}}$$

- For  $Bi < 0.1$  the solid temperature distribution can be considered spatially uniform (depends only on the time):  $T(\mathbf{x}, t) \approx T(t)$ .
  - **The lumped capacitance method provides a solution for  $T(t)$ .**
- For  $Bi \geq 0.1$  the local solid temperatures depend on the position and time.
  - **$T(\mathbf{x}, t)$  solutions to the heat diffusion equation can be evaluated by analytical (exact and approximate) or numerical means.**

# One-Dimensional, Transient Conduction – Gov. Eqs.

Transient conduction can be described in 1D for the case of a plane wall, infinite cylinder and a sphere through the heat equation.



Simplifying assumptions:

- no thermal energy generation; and
- constant thermal conductivity.

## Heat Diffusion Equation

$$\nabla \cdot (k \nabla T) + \dot{q} = \rho c \frac{\partial T}{\partial t}$$

## Plane Wall

$$\frac{\partial^2 T}{\partial x^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

## Infinite Cylinder

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

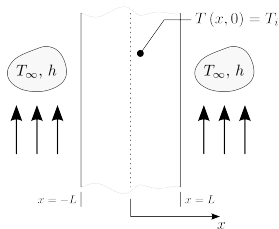
## Sphere

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

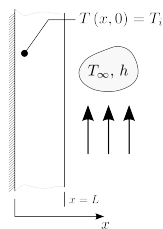
$\alpha = k/(\rho c)$  – Thermal diffusivity

# One-Dimensional, Transient Conduction in a Plane Wall

Symmetrical Convection Conditions



Insulated Surface and Convective Surface



## Governing Equation

$$\frac{\partial^2 T}{\partial x^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

## Initial Condition

$$T(x, t = 0) = T_i$$

## 8 Independent Variables

$$T = f(x, \alpha, t, T_i, k, L, h, T_\infty)$$

## Boundary Conditions

$$\left. \frac{\partial T}{\partial x} \right|_{x=0} = 0$$
$$-k \left. \frac{\partial T}{\partial x} \right|_{x=L} = h[T(L, t) - T_\infty]$$

# One-Dimensional, Transient Conduction in a Plane Wall

Non-dimensionalization:

- $\theta^* = \frac{\theta}{\theta_i} = \frac{T - T_\infty}{T_i - T_\infty}$ 
  - $0 \leq \theta^* \leq 1$
- $x^* = \frac{x}{L}$ 
  - $0 \leq x^* \leq 1$
- $Fo = t^* = \frac{\alpha t}{L^2}$
- $Bi = \frac{hL}{k}$

Governing Equation

$$\frac{\partial^2 \theta^*}{\partial x^{*2}} = \frac{\partial \theta^*}{\partial Fo}$$

Initial Condition

$$\theta^*(x^*, 0) = 1$$

3 Independent Variables

$$\theta^* = f(x^*, Fo, Bi)$$

$\theta^*$  – Dimensionless local temperature difference

$Fo$  – Fourier number

$x^*$  – Dimensionless position

Boundary Conditions

$$\left. \frac{\partial \theta^*}{\partial x^*} \right|_{x^*=0} = 0$$

$$\left. \frac{\partial \theta^*}{\partial x^*} \right|_{x^*=1} = -Bi\theta^*(1, Fo)$$

# One-Dimensional, Transient Conduction in a Plane Wall

## Dimensionless Local Temperature Difference

- **Exact Solution**

The exact solution for the problem is given in the form of an infinite series.

$$\theta^*(x^*, Fo) = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) \cos(\zeta_n x^*)$$

For the geometry under consideration (plane wall),  $C_n$  and  $\zeta_n$  are functions of  $Bi$ .  $C_n$  and  $\zeta_n$  are commonly given in tables.

- **Approximate Solution: One-term Approx. (Valid for  $Fo > 0.2$ )**

$$\theta^*(x^*, Fo) = \frac{T(x^*, Fo) - T_{\infty}}{T_i - T_{\infty}} = \underbrace{C_1 \exp(-\zeta_1^2 Fo)}_{\theta_0^*(Fo) = \theta^*(0, Fo)} \cos(\zeta_1 x^*)$$

$\theta_0^*$  – Midplane ( $x^* = 0$ ) dimensionless temperature difference

# One-Dimensional, Transient Conduction in a Plane Wall

## Dimensionless Mean Temperature Difference

- **Exact Solution**

The exact solution for the problem is given in the form of an infinite series.

$$\bar{\theta}^*(Fo) = \sum_{n=1}^{\infty} \frac{\sin(\zeta_n)}{\zeta_n} C_n \exp(-\zeta_n^2 Fo)$$

- **Bi**  $\rightarrow$  **0** – The exact solution becomes equal to the lumped capacitance method (LCM) solution (considering *Bi* and *Fo* defined with  $L_c = V/A_s$ ):

$$\bar{\theta}^*(Fo) = \theta_{\text{LCM}}^*(Fo) = \exp(-Bi.Fo)$$

- **Approximate Solution: One-term Approx. (Valid for  $Fo > 0.2$ )**

$$\bar{\theta}^*(Fo) = \frac{\bar{T}(Fo) - T_{\infty}}{T_i - T_{\infty}} = \frac{\sin \zeta_1}{\zeta_1} \theta_0^*(Fo)$$



# One-Dimensional, Transient Conduction in a Plane Wall

## Fractional Energy Loss/Gain to/from the Surrounding Fluid

$$\frac{Q(Fo)}{Q_0} = 1 - \bar{\theta}^*(Fo)$$

- $Q(Fo) [= \rho Vc (T_i - \bar{T}(Fo))]$  – Total thermal energy transfer from/to the wall over the time interval  $t [= FoL^2/\alpha]$ .
- $Q_0 [= \rho Vc (T_i - T_\infty)]$  – Initial thermal energy of the wall relative to the fluid temperature, *i.e.*, maximum possible energy transfer from/to the wall if the process continues to time  $t = \infty$ .

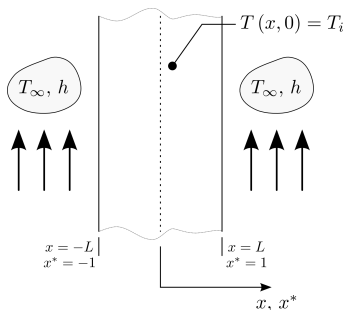
## Boundary Condition at $x^* = 1$ : Constant Surface Temperature

The foregoing solutions for  $\theta^*$ ,  $\bar{\theta}^*$ , and  $Q/Q_0$  are also applicable for a prescribed temperature boundary condition at  $x = L$  ( $T(L, t) = T_s$ ) since this is equivalent to consider  $h = \infty$  ( $Bi = \infty$ ) and  $T_\infty = T_s$ .

# One-Dimensional, Transient Conduction in a Plane Wall

## Heat Removal ( $T_i > T_\infty$ ): Convection Cooling

Numerical and One-Term Approximation Solutions



### 3 Case Studies:

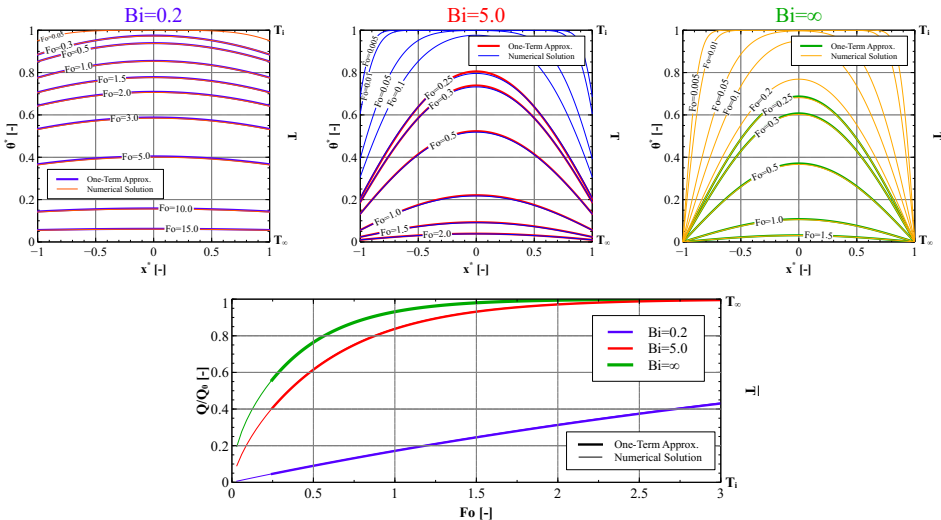
- $Bi = 0.2$ ;
- $Bi = 5.0$ ; and
- $Bi = \infty$ .
  - Negligible convection resistance: equivalent to prescribe a constant surface temperature ( $T_s$ ) equal to  $T_\infty$ .

$$\Delta E_{st} = -Q, \quad Q > 0$$

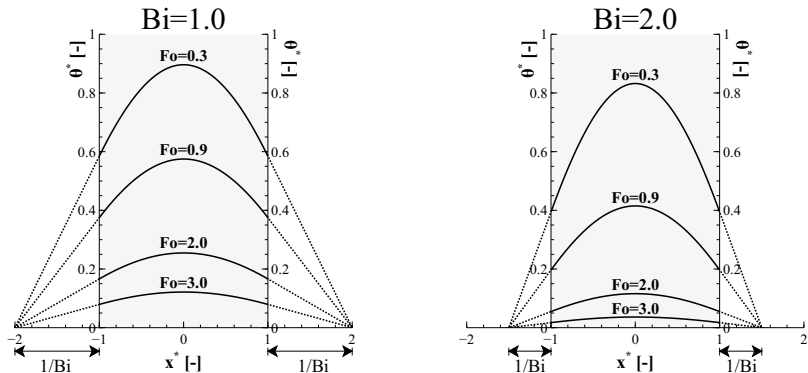
$\Delta E_{st}$  – Change in the thermal energy stored

# One-Dimensional, Transient Conduction in a Plane Wall

## Heat Removal – Numerical and One-Term Approximation Solutions



# One-Dimensional, Transient Conduction in a Plane Wall

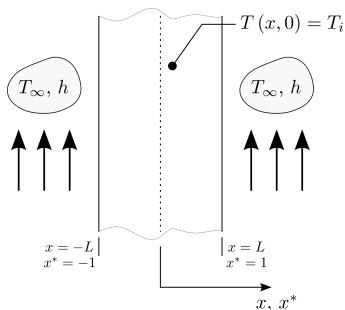


- At any time instant during an unsteady conduction process, the extensions of the tangents to the curves at the points  $x^* = \pm 1$  intersect the axis perpendicular to  $\theta^* = 0$  at the points  $\pm(1 + \frac{1}{Bi})$ .
- This evidence is also observed for long rods and spheres and is due to the mathematical formulation of the convective surface boundary condition.

# One-Dimensional, Transient Conduction in a Plane Wall

## Heat Removal ( $T_i > T_\infty$ ): Convection Cooling

Numerical and One-Term Approximation Solutions



### 3 Case Studies:

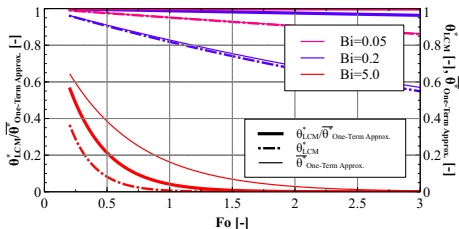
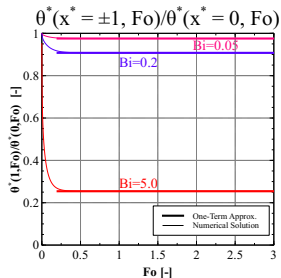
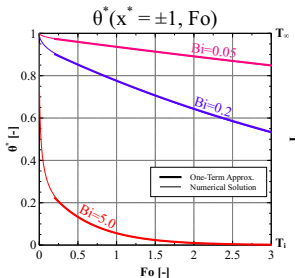
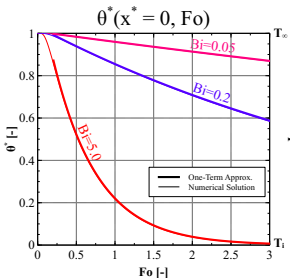
- $Bi = 0.05$ ;
- $Bi = 0.2$ ; and
- $Bi = 5.0$ .

$$\Delta E_{st} = -Q, \quad Q > 0$$

$\Delta E_{st}$  – Change in the thermal energy stored

# One-Dimensional, Transient Conduction in a Plane Wall

## Heat Removal – Numerical and One-Term Approximation Solutions



$$\frac{\theta^*(1, t^*)}{\theta^*(0, t^*)} = f(Bi, Fo).$$

However, as  $Fo \uparrow$   $\frac{\theta^*(1, t^*)}{\theta^*(0, t^*)} \rightarrow \cos(\zeta_1)$

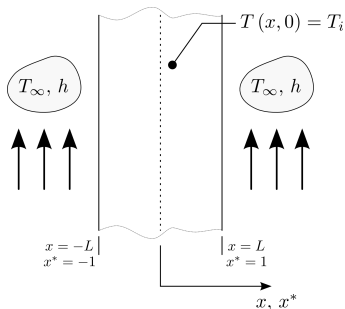
**One-Term Approximation:**

$$\frac{\theta^*(1, t^*)}{\theta^*(0, t^*)} = \begin{cases} 0.976 & Bi = 0.05 \\ 0.908 & Bi = 0.2 \\ 0.254 & Bi = 5.0 \end{cases}$$

# One-Dimensional, Transient Conduction in a Plane Wall

## Heat Addition ( $T_\infty > T_i$ ): Convection Heating

Numerical and One-Term Approximation Solutions



### 3 Case Studies:

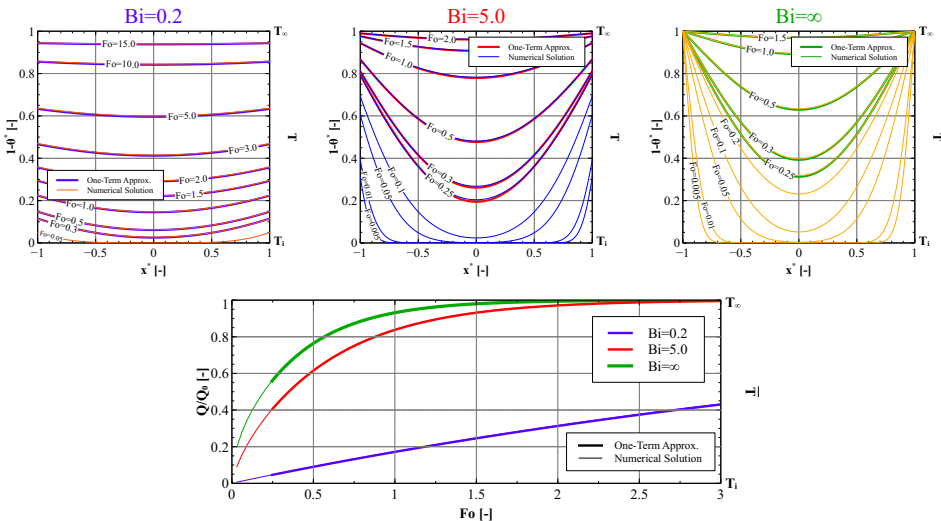
- $Bi = 0.2$ ;
- $Bi = 5.0$ ; and
- $Bi = \infty$ .
  - Negligible convection resistance: equivalent to prescribe a constant surface temperature ( $T_s$ ) equal to  $T_\infty$ .

$$\Delta E_{st} = -Q, \quad Q < 0$$

$\Delta E_{st}$  – Change in the thermal energy stored

# One-Dimensional, Transient Conduction in a Plane Wall

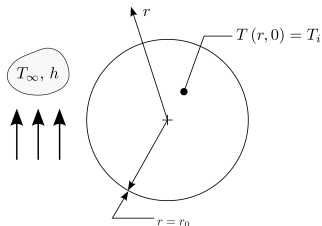
## Heat Addition – Numerical and One-Term Approximation Solutions





# One-Dimen., Transient Conduction in Radial Systems

Infinite Cylinder or Sphere  
Heated/Cooled by Convection



Infinite Cylinder - Gov. Equation

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

Sphere - Governing Equation

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

Boundary Conditions

$$\left. \frac{\partial T}{\partial r} \right|_{r=0} = 0$$

$$-k \left. \frac{\partial T}{\partial r} \right|_{r=r_0} = h [T(r_0, t) - T_\infty]$$

Initial Condition

$$T(r, t = 0) = T_i$$

# One-Dimen., Transient Conduction in Radial Systems

Non-dimensionalization:

- $\theta^* = \frac{\theta}{\theta_i} = \frac{T - T_\infty}{T_i - T_\infty}$ 
  - $0 \leq \theta^* \leq 1$
- $r^* = \frac{r}{r_0}$ 
  - $0 \leq r^* \leq 1$
- $Fo = t^* = \frac{\alpha t}{r_0^2}$
- $Bi = \frac{hr_0}{k}$

Initial Condition

$$\theta^*(r^*, 0) = 1$$

Infinite Cylinder - Gov. Equation

$$\frac{1}{r^*} \frac{\partial}{\partial r^*} \left( r^* \frac{\partial \theta^*}{\partial r^*} \right) = \frac{1}{\alpha} \frac{\partial \theta^*}{\partial Fo}$$

Sphere - Governing Equation

$$\frac{1}{r^{*2}} \frac{\partial}{\partial r^*} \left( r^{*2} \frac{\partial \theta^*}{\partial r^*} \right) = \frac{1}{\alpha} \frac{\partial \theta^*}{\partial Fo}$$

Boundary Conditions

$$\left. \frac{\partial \theta^*}{\partial r^*} \right|_{r^*=0} = 0$$

$$\left. \frac{\partial \theta^*}{\partial r^*} \right|_{r^*=1} = -Bi\theta^*(1, Fo)$$

# One-Dimen., Transient Conduction in Radial Systems

## Dimensionless Local Temperature Difference – Exact Solutions

The exact solutions for the infinite cylinder and sphere are given in the form of infinite series.

### Infinite Cylinder

$$\theta^*(r^*, Fo) = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) J_0(\zeta_n r^*)$$

### Sphere

$$\theta^*(r^*, Fo) = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) \frac{1}{\zeta_n r^*} \sin(\zeta_n r^*)$$

$C_n$  and  $\zeta_n$  are functions of  $Bi$  and the geometry under consideration (long rod or sphere).  $C_n$  and  $\zeta_n$  are commonly given in tables.

# One-Dimen., Transient Conduction in Radial Systems

Approximate Solutions: One-term Approximation (Valid for  $Fo > 0.2$ )

	Infinite Cylinder	Sphere
$\theta^*(r^*, Fo)$	$\theta_0^*(Fo) J_0(\zeta_1 r^*)$	$\theta_0^*(Fo) \frac{1}{\zeta_1 r^*} \sin(\zeta_1 r^*)$
$\theta_0^*(Fo)$	$C_1 \exp(-\zeta_1^2 Fo)$	
$\bar{\theta}^*(Fo)$	$\frac{2J_1(\zeta_1)}{\zeta_1} \theta_0^*(Fo)$	$\frac{3\theta_0^*(Fo)}{\zeta_1^3} [\sin(\zeta_1) - \zeta_1 \cos(\zeta_1)]$
$\frac{Q(Fo)}{Q_0}$	$1 - \bar{\theta}^*(Fo)$	

- $\theta_0^*$  - centerline [centerpoint] dimensionless temperature difference for an infinite cylinder [sphere].

## One-Dimen., Transient Conduction in Radial Systems

### Dimensionless Temperature Difference for $Bi \rightarrow 0$

As  $Bi \rightarrow 0$  the exact solution for  $\theta^*(r^*, Fo)$  becomes equal to the lumped capacitance method solution (considering  $Bi$  and  $Fo$  defined with  $L_c = V/A_s$  –  $L_c$  is equal to  $r_0/2$  and  $r_0/3$  for a long cylinder and sphere, respectively):

$$\theta^*(r^*, Fo) \rightarrow \bar{\theta}^*(Fo) = \exp(-Bi.Fo)$$

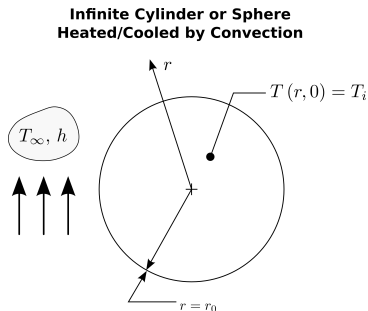
### Boundary Condition at $r^* = 1$ : Constant Surface Temperature

The foregoing solutions for  $\theta^*$ ,  $\bar{\theta}^*$ , and  $Q/Q_0$  are also applicable for a prescribed temperature boundary condition at  $r = r_0$  ( $T(r_0, t) = T_s$ ) since this is equivalent to consider  $h = \infty$  ( $Bi = \infty$ ) and  $T_\infty = T_s$ .

# One-Dimen., Transient Conduction in Radial Systems

## Heat Removal ( $T_i > T_\infty$ ): Convection Cooling

Numerical and One-Term Approximation Solutions



### 3 Case Studies:

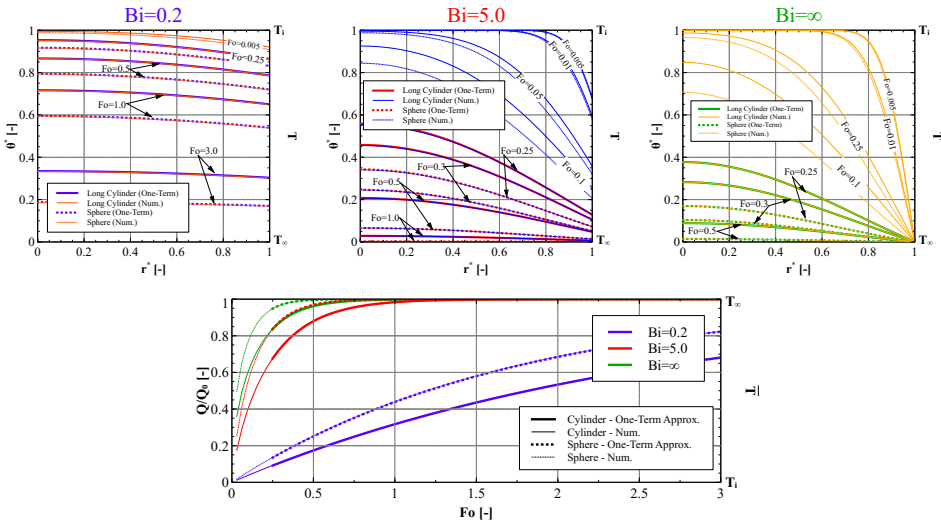
- $Bi = 0.2$ ;
- $Bi = 5.0$ ; and
- $Bi = \infty$ .
  - Negligible convection resistance: equivalent to prescribe a constant surface temperature ( $T_s$ ) equal to  $T_\infty$

$$\Delta E_{st} = -Q, \quad Q > 0$$

$\Delta E_{st}$  – Change in the thermal energy stored

# One-Dimen., Transient Conduction in Radial Systems

## Heat Removal – Numerical and One-Term Approximation Solutions



## Final Remarks (1/2)

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- Biot number provides an estimation for the relevance of temperature spatial gradients in a heat conduction process within a solid concurrent with convection across its boundaries.

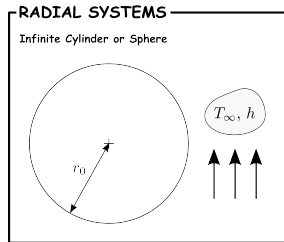
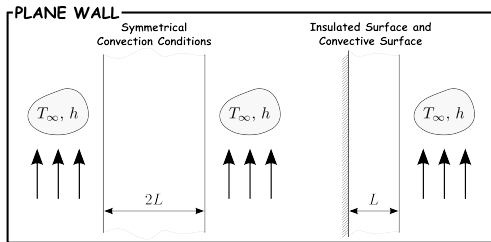
For a one-dimensional, transient heat conduction process if:

- $Bi < 0.1$ : the spatial gradients are not relevant; consequently, the lumped capacitance method can be applied;
  - $Bi \geq 0.1$ : the spatial gradients are relevant; consequently, the one-term approximation to the exact solution – particularly recommended for  $Fo > 0.2$  – or a numerical procedure should be applied to evaluate the temporal and spatial solid temperature distribution profiles.
- The one-term approx. for  $Fo > 0.2$  results in an error below 2%.
  - Heisler/Gröber charts (transient temperature and heat transfer charts) provide a graphical representation for  $\theta_0^*$ ,  $\theta^*/\theta_0^*$ , and  $Q/Q_0$  obtained with the single-term approximation of the exact solution.



## Final Remarks (2/2) – $L_c$ for Biot and Fourier Numbers

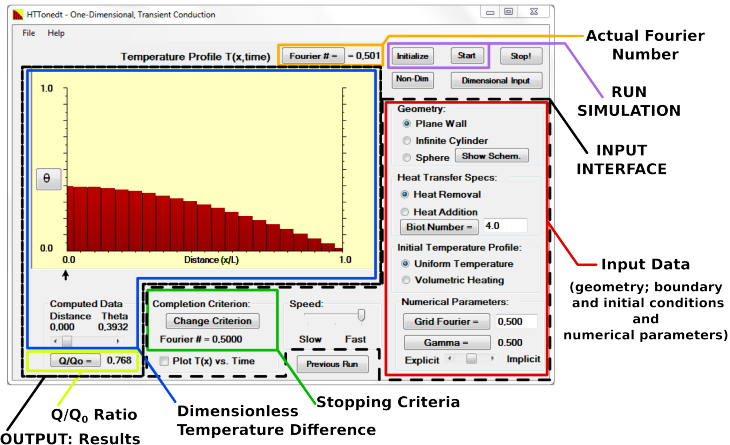
	$L_c$ – Characteristic length <sup>1</sup>		
	Plane Wall	Inf. Cylinder	Sphere
Conservative $Bi$ Criterion (relevance of temp. spatial gradients)	$L$	$r_0$	$r_0$
Lumped capacitance method – $L_c = V/A_s$	$L$	$r_0/2$	$r_0/3$
Analytical and numerical solutions for $\theta^*(x^*, Fo)$	$L$	$r_0$	$r_0$



<sup>1</sup>Find the  $L_c$  value ( $L$  and  $r_0$ ) in accordance with the accompanying figure.

# Exploring the Software Module (1/4)

## Software module – HTTonedt.exe (Version 5.0.0.2)



- The module solves the one-dimensional, transient heat equation through a finite-volume approach for a plane wall, infinite cylinder, and sphere.

# Exploring the Software Module (2/4)

## Software module – HTTonedt.exe (Version 5.0.0.2)

The screenshot displays the HTTonedt software interface for "One-Dimensional, Transient Conduction". It features two dialog boxes for stopping criteria, a main simulation window, and a results panel.

- Two Stopping Criteria:** Two dialog boxes are shown. The left one is titled "Specified Fourier Number" and has "Fourier # = 0.5" entered. The right one is titled "Specified Dim. Local Temp. Diff." and has "Specified Theta" selected with "Distance = 0.000" and "Theta = 500" entered.
- Actual Fourier Number:** A label points to the "Fourier # = 0.5" field in the left dialog.
- RUN SIMULATION:** A label points to the "Enter/Return" button in the right dialog.
- INPUT INTERFACE:** A label points to the main simulation window.
- Input Data (geometry; boundary and initial conditions and numerical parameters):** A label points to the right panel, which includes "Initial Temperature Profile" (Uniform Temperature selected), "Numerical Parameters" (Grid Fourier = 0.500, Gamma = 0.500), and "Explicit" vs "Implicit" options.
- Stopping Criteria:** A label points to the "Change Criterion" button in the main window.
- Dimensionless Temperature Difference:** A label points to the "Theta" value in the "Computed Data" table.
- Q/Q<sub>0</sub> Ratio OUTPUT: Results:** A label points to the "Q/Q<sub>0</sub> = 0.768" value in the main window.

Distance (x/L)	Theta
0.000	0.3932

Computed Data: Q/Q<sub>0</sub> = 0.768

- The module ends the simulation for two possible criteria: (a) specified Fourier number; and (b) specified dimensional local temperature difference.

# Exploring the Software Module (3/4)

## Completion Criteria

The module terminates the simulation for two possible criteria:

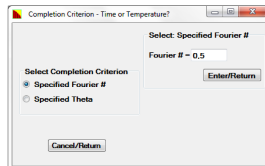
### 1. Specified Fourier number – $Fo$ ; and

- For evaluation of the temperature distribution profiles and the ratio  $Q/Q_0$  at a specific time instant.

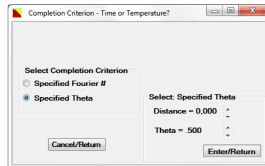
### 2. Specified dimensionless local temperature difference – $\theta^*(x^*, Fo)$ .

- For the evaluation of the elapsed time, temperature distribution profiles, and the ratio  $Q/Q_0$ .

#### 1. Specified Fourier Number



#### 2. Specified Dimensionless Local Temperature Difference



## Exploring the Software Module (4/4)

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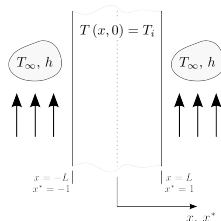
### Spatial Discretization

- Two spatial discretization schemes (meshes) are available in the software module.
- The finest mesh has two times the cell count of the default mesh and, consequently, should provide more accurate results; however, at the expense of longer computation times.
- The finest grid is taken into account (activated) for the calculations once the default window size is changed.
- To revert to the default mesh, the user should restart the software module.
- The module application examples that follow (next slides) consider the default mesh.

# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Problem Statement

Consider a 0.1 m ( $2L$ ) thick plane wall initially at  $T_i = 180\text{ }^\circ\text{C}$  that is suddenly cooled with a fluid at  $T_\infty = 20\text{ }^\circ\text{C}$  and with  $h = 2200\text{ W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$ . The wall material has a thermal conductivity ( $k$ ), density ( $\rho$ ), and specific heat ( $c$ ) equal to  $110\text{ W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ ,  $8530\text{ kg}\cdot\text{m}^{-3}$ , and  $380\text{ J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$ , respectively.



After 40 s of cooling, evaluate the following using the module:

1. temperature distribution profile,  $T(-L \leq x \leq L)$ ;
2. fractional energy loss,  $Q/Q_0$ ;
3. average wall temperature,  $\bar{T}$ ; and
4. compare the average wall temperature computed with the module with the temperature predicted by the lumped capacitance method.

# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Module Application

### Preliminary Calculations

<u>Biot Number</u>	<u>Thermal Diffusivity</u>	<u>Fourier Number</u>
$Bi = 1.00$	$\alpha = 3.39 \times 10^{-5} \text{ m}^2 \cdot \text{s}^{-1}$	$Fo(t = 40 \text{ s}) \approx 0.54$

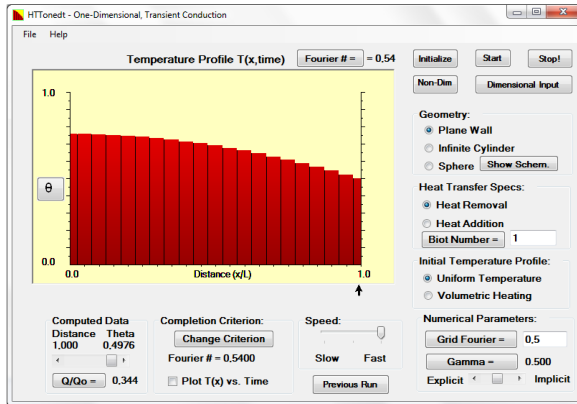
### Module Input Data

<u>1 - Geometry:</u> "Plane Wall"	<u>2 - Heat Transfer Specs:</u> "Heat Removal" "Biot Number = 1"	<u>3 - Initial Temperature Profile:</u> "Uniform Temperature"
<u>4 - Numerical Parameters:</u> "Grid Fourier = 0,5" "Gamma = 0.500" "Default Mesh"		<u>5 - Completion Criteria:</u> "Fourier # = 0,5400"

# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Module Application and Results

### Module Results





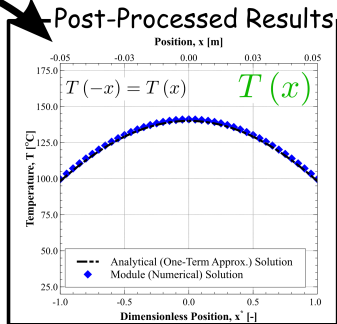
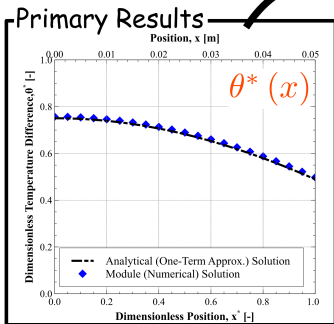
# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Results Analysis (1/3)

### 1. Temperature distribution profile, $T(-L \leq x \leq L)$

$$\theta^*(x) = \frac{T(x) - T_\infty}{T_i - T_\infty}$$

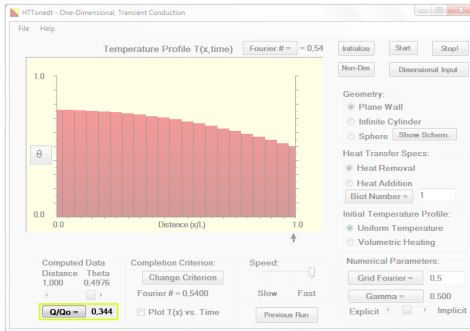
$$\begin{aligned} T_i &= 180^\circ\text{C} \\ T_\infty &= 20^\circ\text{C} \end{aligned}$$



# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Results Analysis (2/3)

### 2. Fractional energy loss, $Q/Q_0$



$$Q/Q_0 = 0.344 \quad \left( (Q/Q_0)_{\text{Analytic Sol.}} \approx 0.339 \right)$$

# Exploring the Module - Cooling a Plane Wall

## Module Application Example I: Results Analysis (3/3)

### 3. Average wall temperature, $\bar{T}$

$$\begin{aligned}\bar{\theta}^* &= \int_0^1 \theta^*(x^*) dx^* = \\ &= 1 - Q/Q_0 = 0.656\end{aligned}$$

$$\begin{aligned}\bar{\theta}^* &= \frac{\bar{T} - T_\infty}{T_i - T_\infty} \Leftrightarrow \\ \Leftrightarrow &\boxed{\bar{T} = 124.96^\circ\text{C}}\end{aligned}$$

$$(\bar{T}_{\text{Analytic Sol.}} \approx 125.79^\circ\text{C})$$

### 4. Average wall temperature – module vs. lumped capacitance method results

$$\begin{aligned}\theta_{\text{LCM}}^* &= \exp(-Bi \cdot Fo) = \exp\left(-\frac{h\alpha t}{kL}\right) = \\ &= \exp\left(-\frac{2200 \times 3.39 \times 10^{-5} \times 40}{110 \times 0.05}\right) \Leftrightarrow \\ &\Leftrightarrow \theta_{\text{LCM}}^* \approx 0.581\end{aligned}$$

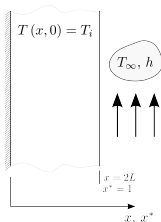
$$\theta_{\text{LCM}}^* = \frac{T_{\text{LCM}} - T_\infty}{T_i - T_\infty} \Leftrightarrow \boxed{T_{\text{LCM}} = 112.960^\circ\text{C}}$$

A relative **error of about 10%** is observed for  $T_{\text{LCM}}$  in relation to the module solution ( $\bar{T}_{\text{Mod.}} = 124.96^\circ\text{C}$ ).

# Exploring the Module - Heating a Plane Wall

## Module Appl. Example II: Problem Statement

Consider the same plane wall of Example I (same thermophysical properties and geometrical parameters) initially at  $T_i = 20^\circ\text{C}$ . One surface is perfectly insulated while the other is suddenly exposed to a fluid at  $T_\infty = 180^\circ\text{C}$  and with  $h = 2200\text{ W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$ .



Evaluate the following using the module:

1. elapsed time,  $t$ , to observe a temperature equal to  $100^\circ\text{C}$  at the insulated surface (*i.e.*,  $T(x = 0, t) = 100^\circ\text{C}$ );
2. temperature at  $x = 0.08\text{ m}$  and at the time instant of 1.; and
3. thermal energy absorbed per unit active surface area,  $Q/A_s$ , at the time instant of 1.

# Exploring the Module - Heating a Plane Wall

## Module Appl. Example II: Module Application

### Preliminary Calculations

<u>Biot Number</u>	<u>Thermal Diffusivity</u>	<u>Dim. Local Temp. Diff.</u>
$Bi = 2.00$	$\alpha = 3.39 \times 10^{-5} \text{ m}^2 \cdot \text{s}^{-1}$	$\theta^*(0, Fo) = 0.5$

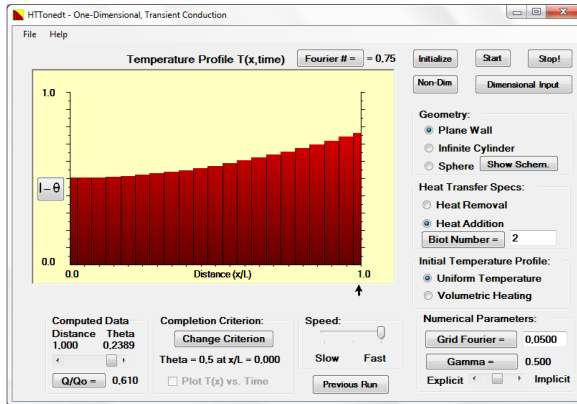
### Module Input Data

<u>1 - Geometry:</u> "Plane Wall"	<u>2 - Heat Transfer Specs:</u> "Heat Addition" "Biot Number = 2"	<u>3 - Initial Temperature Profile:</u> "Uniform Temperature"
<u>4 - Numerical Parameters:</u> "Grid Fourier = 0,05" "Gamma = 0.500" "Default Mesh"		<u>5 - Completion Criteria:</u> "Theta=0.5 at x/L=0.000"

# Exploring the Module - Heating a Plane Wall

## Module Appl. Example II: Module Application and Results

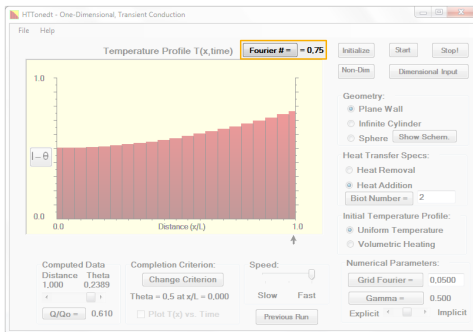
### Module Results



# Exploring the Module - Heating a Plane Wall

## Module Application Example II: Results Analysis (1/3)

1. Elapsed time to observe  $T(x = 0, t) = 100^\circ\text{C}$



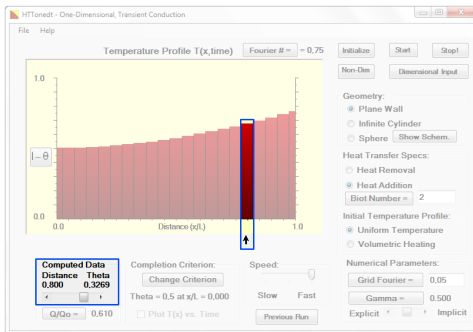
$$Fo = 0.75 \Leftrightarrow t = \frac{0.1^2}{3.39 \times 10^{-5}} \times 0.75 \Leftrightarrow t \approx 3 \text{ min. and } 41 \text{ sec.}$$

$$((\theta^*(0, Fo))_{\text{Analytic Sol.}} = 0.5 \Rightarrow t_{\text{Analytic Sol.}} \approx 3 \text{ min. and } 38 \text{ sec.})$$

# Exploring the Module - Heating a Plane Wall

## Module Application Example II: Results Analysis (2/3)

2. Temperature at  $x = 0.08$  m when  $T(x = 0, t) = 100$  °C



$$\theta^* (x^* = 0.8, Fo) = 0.3269 \Rightarrow T(x = 0.08 \text{ m}, t) \approx 127.7 \text{ }^\circ\text{C}$$

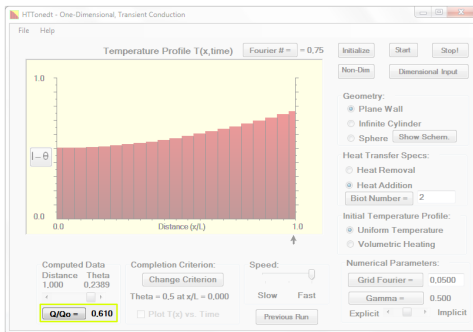
$$\left( (T(x = 0.08 \text{ m}, t)) \right)_{\text{Analytic Sol.}} \approx 127.9 \text{ }^\circ\text{C}$$



# Exploring the Module - Heating a Plane Wall

## Module Application Example II: Results Analysis (3/3)

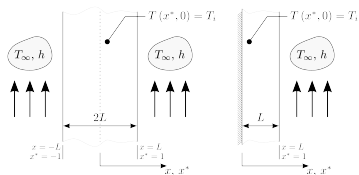
3. Absorbed energy per unit active surface area,  $Q/A_s$ , when  $T(x = 0, t) = 100\text{ }^\circ\text{C}$



$$Q/A_s = \rho L_c c (Q/Q_0) \theta_i = 8530 \times 0.1 \times 380 \times 0.610 \times (20 - 180) \Leftrightarrow$$
$$\Leftrightarrow Q/A_s \approx -3.16 \times 10^7 \text{ J} \cdot \text{m}^{-2} \quad \left( (Q/A_s)_{\text{Analytic Sol.}} \approx -3.07 \times 10^7 \text{ J} \cdot \text{m}^{-2} \right)$$

# Useful Relations

## Plane Wall



$$\theta^* = \frac{T(x^*, Fo) - T_\infty}{T(x^*, 0) - T_\infty}$$

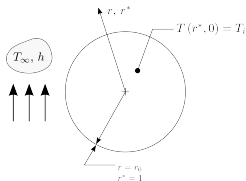
$$x^* = \frac{x}{L}$$

$$Fo = \frac{\alpha t}{L^2}$$

$$Bi = \frac{hL}{k}$$

$$\alpha = k / (\rho c)$$

## Radial Systems



$$\theta^* = \frac{T(r^*, Fo) - T_\infty}{T(r^*, 0) - T_\infty}$$

$$r^* = \frac{r}{r_0}$$

$$Fo = \frac{\alpha t}{r_0^2}$$

$$Bi = \frac{hr_0}{k}$$

$$Q_0 = \rho V c (T_i - T_\infty)$$